

## Kinetic analysis of open- and closed-state inactivation transitions in human Kv4.2 A-type potassium channels

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1. We studied the gating kinetics of Kv4.2 channels, the molecular substrate of neuronal somatodendritic A-type currents. For this purpose wild-type and mutant channels were transiently expressed in the human embryonic kidney (HEK) 293 cell line and currents were measured in the whole-cell patch-clamp configuration.
2. Kv4.2 channels inactivated from pre-open closed state(s) with a mean time constant of 959 ms at  $-50$  mV. This closed-state inactivation was not affected by a deletion of the Kv4.2 N-terminus ( $\Delta 2-40$ ).
3. Kv4.2 currents at  $+40$  mV inactivated with triple-exponential kinetics. A fast component ( $\tau = 11$  ms) accounted for 73%, an intermediate component ( $\tau = 50$  ms) for 23% and a slow component ( $\tau = 668$  ms) for 4% of the total decay.
4. Both the fast and the intermediate components of inactivation were slowed by a deletion of the Kv4.2 N-terminus ( $\tau = 35$  and 111 ms) and accounted for 33 and 56%, respectively, of the total decay. The slow component was moderately accelerated by the truncation ( $\tau = 346$  ms) and accounted for 11% of the total Kv4.2 current inactivation.
5. Recovery from open-state inactivation and recovery from closed-state inactivation occurred with similar kinetics in a strongly voltage-dependent manner. Neither recovery reaction was affected by the N-terminal truncation.
6. Kv4.2  $\Delta 2-40$  channels displayed slowed deactivation kinetics, suggesting that the N-terminal truncation leads to a stabilization of the open state.
7. Simulations with an allosteric model of inactivation, supported by the experimental data, suggested that, in response to membrane depolarization, Kv4.2 channels accumulate in the closed-inactivated state(s), from which they directly recover, bypassing the open state.

Voltage-activated potassium (Kv) channels with rapid (A-type) inactivation play important roles in pre- and postsynaptic neuronal excitation (Pongs, 1999). A-type Kv channels related to the *Shal* (Kv4) gene family (Baldwin *et al.* 1991; Pak *et al.* 1991; Serodio *et al.* 1994, 1996) mediate both the transient outward current ( $I_{to}$ ) in cardiac myocytes (Dixon *et al.* 1996) and the neuronal subthreshold A-type current (Serodio *et al.* 1994). The Kv4 subfamily has four members, Kv4.1, Kv4.2 and two splice variants of Kv4.3, cloned from rat (Blair *et al.* 1991; Dixon & McKinnon, 1994; Serodio *et al.* 1996; Ohya *et al.* 1997; Tsaour *et al.* 1997), mouse (Pak *et al.* 1991; Serodio *et al.* 1994), dog (Dixon *et al.* 1996), or human cDNA libraries (Kong *et al.* 1998; Dilks *et al.* 1999; Zhu *et al.* 1999; Isbrandt *et al.* 2000).

Immunocytochemical studies have shown that the subcellular distribution of neuronal rat Kv4.2 channels is

restricted to the somatodendritic area (Sheng *et al.* 1992), and the high abundance of Kv4.2 in the soma and dendrites led to the hypothesis that these channels may have an important influence on postsynaptic neuronal signal transduction. The discovery of Kv4.2 channel protein clusters in the postsynaptic membrane supported this view (Alonso & Widmer, 1997). Many important postsynaptic functions have been demonstrated for dendritic A-type channels. In hippocampal pyramidal neurons, where the A-type channel density increases with distance from the soma, activation of Kv4.2 channels may prevent the back-propagation of action potentials (Hoffman *et al.* 1997). In addition, the rapid activation of A-type channels may protect the postsynaptic membrane from excessive depolarization. On the other hand, inactivation of Kv4.2 channels by subthreshold EPSPs

can lead to spike amplification, which provides a possible explanation for the Hebbian associativity found in distal dendrites (Magee & Johnston, 1997). More recently, Schoppa & Westbrook (1999) showed that A-type channels localized to the dendritic spines of GABAergic interneurons in the olfactory bulb are necessary to counterbalance fast glutamatergic EPSPs, thereby allowing a prolonged inhibitory synaptic transmission in a local feedback loop. The inhibitory synaptic transmission becomes shorter and stronger when the A-type channels are inactivated by subthreshold membrane depolarization.

The gating properties of somatodendritic A-type channels reflect their functional roles. Not only do these channels rapidly inactivate close to the action potential firing threshold, but they also recover from inactivation in the millisecond time range. This property distinguishes somatodendritic from other A-type channels. For example, *Shaker*-related A-type channels like Kv1.4 slowly recover from inactivation within a time range of many seconds. The inactivation mechanism of *Shaker* channels is well understood. It can be described as a 'ball-and-chain' reaction (Hoshi *et al.* 1990, 1991) where a tethered N-terminal inactivation domain of the Kv $\alpha$  subunit (Zagotta *et al.* 1990), or of an accessory Kv $\beta$  subunit (Rettig *et al.* 1994), binds to a ball receptor near the inner entrance of the pore (Isacoff *et al.* 1991) thereby preventing ion flux through the open channel. The fast N-type inactivation of *Shaker* channels is followed by a slower C-type inactivation (Hoshi *et al.* 1991), from which the channels recover slowly. By contrast, Jerng and coworkers (Jerng & Covarrubias, 1997; Jerng *et al.* 1999) showed that the inactivation of Kv4.1 channels, when expressed *in vitro*, does not require channel opening. The channels readily enter inactivated state(s) from pre-open closed state(s). In fact, this pathway appears to be the only one that leads to cumulative Kv4.1 channel inactivation; i.e. open channels have to close before they accumulate in the inactivated state(s). In the Kv4.1 gating model proposed by Jerng *et al.* (1999) this cumulative inactivation involves two sequential states, a closed-inactivated state and a deep-inactivated state.

In order to gain further insight into the gating behaviour of somatodendritic A-type channels, we investigated the kinetics of currents mediated by wild-type and N-terminally truncated Kv4.2 channels expressed in human embryonic kidney (HEK) 293 cells. We performed a detailed kinetic analysis of both closed- and open-state inactivation and recovery from inactivation.

Our results indicate that the Kv4.1 gating model of Jerng *et al.* (1999) cannot be used to simulate A-type currents mediated by Kv4.2 channels. In particular, the fast recovery from inactivation, a hallmark of somatodendritic A-type channels, which has been observed in our experiments, could not be simulated. Our experimental data could be most accurately reproduced by simulations based on an allosteric model of inactivation, where all

closed states are connected to inactivated states, but not to a deep-inactivated state. Simulations with this model yielded fast recovery kinetics after accumulation in the closed-inactivated state(s), just as observed experimentally for Kv4.2 channels.

## METHODS

### Cell culture and heterologous channel expression

Both wild-type human Kv4.2 channels (EMBL data bank accession no. AJX 010969) and an N-terminal deletion mutant, which lacks amino acid residues 2–40 ( $\Delta 2-40$ ), were transiently overexpressed in human embryonic kidney (HEK) 293 cells, as described previously (Zhu *et al.* 1999). Constructs of human Kv4.2 channels with a larger N-terminal truncation (e.g.  $\Delta 2-52$ ) are non-functional (Zhu *et al.* 1999). The cells were grown according to standard culture protocols and plated on poly-L-lysine ( $50 \mu\text{g ml}^{-1}$ )-coated 35 mm plastic dishes at densities between  $1 \times 10^4$  and  $1.5 \times 10^4$  cells per dish. After 7–8 h the cells were transfected using the calcium phosphate technique (Chen & Okayama, 1987) with  $1 \mu\text{g}$  per dish of recombinant Kv4.2 cDNA. We cotransfected cells with  $0.3 \mu\text{g}$  per dish of cDNA for an enhanced version (S56T mutant) of the green fluorescent protein as a reporter gene to identify transfected cells. The precipitate was removed 12 h later by rinsing with phosphate-buffered saline and the cells were used for electrophysiological recordings on the next 2 days.

### Solutions and recording techniques

For electrophysiological recordings, cells were bathed in an external solution containing (mM): NaCl 135, KCl 5, CaCl<sub>2</sub> 2, MgCl<sub>2</sub> 2, Hepes 5 and sucrose 10, with  $0.01 \text{ mg ml}^{-1}$  phenol red; pH 7.4 (NaOH). Recording pipettes, pulled from thin-walled borosilicate glass using a DMZ puller (Zeitz, Augsburg, Germany) and heat polished, had bath resistances between 2.5 and  $3 \text{ M}\Omega$  when filled with an internal solution containing (mM): KCl 115, KF 10, CaCl<sub>2</sub> 1, MgCl<sub>2</sub> 1, EGTA 11, Hepes 10, glutathione 2 and K<sub>2</sub>-ATP 2; pH 7.2 (KOH). All experiments were conducted in the whole-cell configuration of the patch-clamp technique. In some experiments, external Na<sup>+</sup> was replaced by K<sup>+</sup> (130 mM; symmetrical K<sup>+</sup>), or alternatively, both internal K<sup>+</sup> and external Na<sup>+</sup> were replaced by Rb<sup>+</sup> (130 mM; symmetrical Rb<sup>+</sup>). The external solutions with different ionic compositions were applied to the cell surface using a valve-controlled fast local superfusion system (Auto Mate Scientific, San Francisco, CA, USA) driven by a peristaltic pump. All experiments were conducted at room temperature (20–22°C).

### Data acquisition and analysis

Currents were evoked by voltage-jump protocols and recorded using an EPC9 patch-clamp amplifier combined with PULSE software (HEKA Elektronik, Landbrecht, Germany). Signals were filtered at 0.2–4 kHz with low-pass Bessel characteristics, amplified as required, and digitized at 50–1000  $\mu\text{s}$  sample intervals. Capacitive current transients were compensated on-line, and the *P/n* method incorporated in PULSE was used for leak current subtraction (leak holding,  $-100 \text{ mV}$ ; leak size, 20%). Data were stored on a 7600/132 Power Macintosh computer, and the program package PULSEFIT/PULSETOOLS (HEKA Elektronik) was used to analyse current traces. The sigmoidal rising phase of Kv4.2-mediated currents was fitted using a Hodgkin-Huxley-related equation ( $m^4$ ) incorporated in PULSEFIT (Roeper *et al.* 1997). Kv4.2 current inactivation kinetics were fitted with triple-exponential functions and current relaxations caused by channel deactivation with single-exponential functions. The obtained data were further processed using Igor Pro (WaveMetrics, Lake Oswego, OR, USA) and Kaleidagraph (Synergy Software, Reading, PA, USA). Conductance values ( $G$ ) at a given test

**Table 1. Values of simulation parameters used in Model 1**

Parameter	Values
$\alpha_0$ (s <sup>-1</sup> )	133
$z_\alpha$	1.0
$\beta_0$ (s <sup>-1</sup> )	4
$z_\beta$	1.58
$k_{co}$ (s <sup>-1</sup> )	400
$k_{oc}$ (s <sup>-1</sup> )	1100
$k_{oi}$ (s <sup>-1</sup> )	100
$k_{io}$ (s <sup>-1</sup> )	10
$k_{ci}$ (s <sup>-1</sup> )	30
$k_{ic}$ (s <sup>-1</sup> )	5
$k_{id}$ (s <sup>-1</sup> )	8
$k_{di}$ (s <sup>-1</sup> )	0.5

In the models the voltage-dependent rates are of the form:

$$\alpha = \alpha_0 \exp[z_\alpha V/(RT/F)]$$

for the forward rates and:

$$\beta = \beta_0 \exp[-z_\beta V/(RT/F)]$$

for the backward rates, where  $\alpha_0$  and  $\beta_0$  are the rates at 0 mV and  $z_\alpha$  and  $z_\beta$  are the equivalent charges moving up to the transition state. Channel opening and closing are defined by  $k_{co}$  and  $k_{oc}$ , respectively, open-state inactivation by  $k_{oi}$  and  $k_{io}$ , and closed-state inactivation by  $k_{ci}$  and  $k_{ic}$ . Trapping in the deep-inactivated state is defined by  $k_{id}$  and  $k_{di}$ . For Model 1, see Fig. 10.

potential ( $V_m$ ) were calculated from the measured peak amplitude ( $I$ ) and the mean reversal potential for K<sup>+</sup> currents ( $V_{rev} = -79.9 \pm 1.4$  mV;  $n = 5$ ) using the equation  $G = I/(V_m - V_{rev})$ . The voltage dependence of activation and steady-state inactivation of the expressed Kv4.2 channels were analysed with Boltzmann functions of the form  $G/G_{max} = 1/\{1 + \exp[(V_m - V_{1/2})/k]\}^n$ , where the fraction of the maximal conductance,  $G/G_{max}$ , at a given membrane potential,  $V_m$ , depends on the potential for half-maximal activation or inactivation,  $V_{1/2}$ , and the slope factor,  $k$ . Steady-state inactivation data were fitted with a first-order Boltzmann function ( $n = 1$ ). The voltage dependence of activation was fitted with both a first-order ( $n = 1$ ) and a fourth-order Boltzmann function ( $n = 4$ ). Statistical analyses with Student's  $t$  test were conducted using GraphPad Instat (GraphPad Software, San Diego, CA, USA), and pooled data are presented as means  $\pm$  S.E.M.

### Gating model and simulations

For the development of a computational model for the gating of Kv4.2 channels a custom-made program was implemented using the Borland C++ version 4.0 compiler with Intel Pentium instruction set optimization and run on a Pentium computer. We expressed the Markov-state kinetic model as a system of differential equations with voltage-dependent transitions along the activation pathway leading from C<sub>R</sub> (resting state) to C<sub>i</sub> (final pre-open closed state; see Figs 10A and 11A). The voltage-dependent rates are of the form  $\alpha = \alpha_0 \exp[z_\alpha V/(RT/F)]$  for the forward rates and  $\beta = \beta_0 \exp[-z_\beta V/(RT/F)]$  for the backward rates, where  $\alpha_0$  and  $\beta_0$  are the rates at 0 mV and  $z_\alpha$  and  $z_\beta$  are the equivalent charges moving up to the transition state. Numerical time integration was performed using eigenvalue–eigenvector decomposition of the Markov model vector field matrix. Initial conditions for each voltage-clamp sequence computation were computed by direct solving of the steady-state equations for the Markov chain, similar to the effect of a very long clamp pulse at the holding potential. We first used the gating scheme proposed by Jerng *et al.* (1999), which was then modified, and model parameters were allowed to vary, in order to simulate our experimental records. Parameters used in the simulations are given in Tables 1 and 2.

**Table 2. Values of simulation parameters used in Model 2**

Parameter	Values
$\alpha_0$ (s <sup>-1</sup> )	200
$z_\alpha$	1.0
$\beta_0$ (s <sup>-1</sup> )	4
$z_\beta$	1.58
$k_{co}$ (s <sup>-1</sup> )	400
$k_{oc}$ (s <sup>-1</sup> )	1100
$k_{oi}$ (s <sup>-1</sup> )	300
$k_{io}$ (s <sup>-1</sup> )	30
$k_{ci}$ (s <sup>-1</sup> )	30
$k_{ic}$ (s <sup>-1</sup> )	0.1
$f$	0.3
$k_{is}$ (s <sup>-1</sup> )	30
$k_{si}$ (s <sup>-1</sup> )	15

In the allosteric model of Kv4.2 inactivation gating, all closed states (C) are connected to inactivated states via  $k_{ci}$  and  $k_{ic}$ . The coupling of activation and inactivation is defined by an allosteric factor  $f$  (see Fig. 11A). Furthermore, a transition leading from I<sub>O</sub> to I<sub>S</sub> has been added in Model 2 (Fig. 11), which is defined by  $k_{is}$  and  $k_{si}$ . The voltage-dependent reactions along the activation pathway are the same as for Model 1 (see Table 1) with the exception of the  $\alpha_0$  term, which was shifted from 133 s<sup>-1</sup> in Model 1 to 200 s<sup>-1</sup> in Model 2.

## RESULTS

### Currents mediated by Kv4.2 channels

A standard voltage protocol activated the currents mediated by Kv4.2 channels in HEK 293 cells with a 2.5 s pulse to +40 mV following a 2 s hyperpolarizing prepulse to -100 mV. The holding potential during the time between such voltage jumps (15 s) was -80 mV. A second voltage pulse to +40 mV was preceded by a 2 s prepulse to -35 mV to activate endogenous HEK 293 currents (Yu & Kerchner, 1998), if present, in the absence of Kv4.2 currents. The endogenous currents were subtracted from the total K<sup>+</sup> currents if they exceeded 5% of the total current peak amplitude. The results, which reflected exclusively Kv4.2 channel kinetics, are shown in Fig. 1A. Kv4.2 channels displayed strong inactivation with three distinct components (Fig. 1A and B). For wild-type channels the fast component of inactivation ( $\tau_1$ ) measured at +40 mV was  $11.0 \pm 0.7$  ms ( $n = 13$ ) and accounted for  $73.2 \pm 1.7\%$  of the total decay. The fast component was followed by an intermediate and a very slow component of inactivation (Fig. 1B). The intermediate component ( $\tau_2 = 50.1 \pm 2.4$  ms,  $n = 13$ ) accounted for  $22.5 \pm 1.5\%$  and the slow component ( $\tau_3 = 668 \pm 35$  ms,  $n = 13$ ) for  $4.3 \pm 0.3\%$  of the total decay. When we expressed mutant channels lacking 40 amino acids of the N-terminus ( $\Delta 2-40$ ) the resulting currents at +40 mV still exhibited decay kinetics that could be best described by the sum of three exponential terms (Fig. 1A and B). However, the time constants and their relative weights were altered in a characteristic manner. Both the fast and the intermediate components of inactivation were slowed ( $\tau_1 = 35.1 \pm 4.2$  ms and  $\tau_2 = 111 \pm 4$  ms,  $n = 9$ ) and now accounted for

$33.3 \pm 3.1$  and  $55.5 \pm 2.8\%$ , respectively, of the total decay. By contrast, the slow component of Kv4.2  $\Delta 2-40$  inactivation was slightly faster than observed for wild-type channels ( $\tau_3 = 346 \pm 14$  ms,  $n = 9$ ;  $P < 0.01$ ) and accounted for the remaining  $11.2 \pm 2.2\%$  of the total decay. These results are summarized in Fig. 1C. The slowing of the major component of outward current decay suggests a decrease in the rate of Kv4.2 channel open-state inactivation by the  $\Delta 2-40$  mutation. Our findings clearly separate Kv4.2 from *Shaker* B channels, where a similar N-terminal deletion completely abolishes fast inactivation (Hoshi *et al.* 1990).

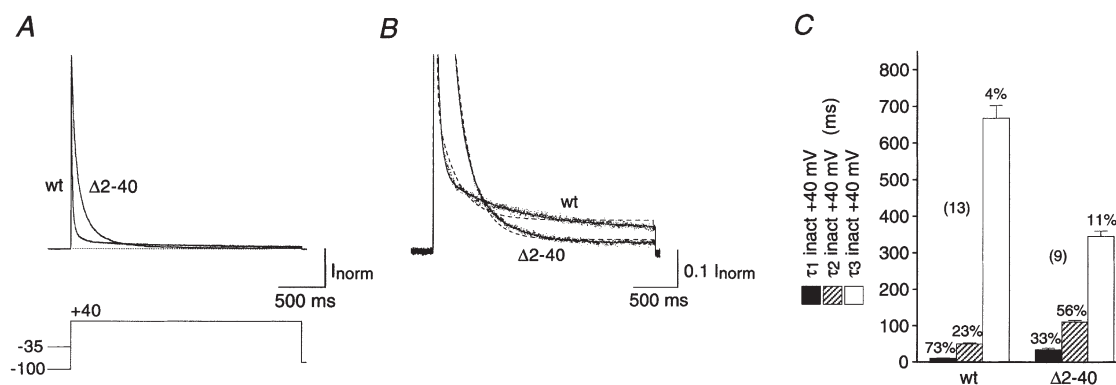
### Activation parameters of Kv4.2 channels

A series of test pulses yielded current families as shown in Fig. 2A. With these voltage-jump protocols we never observed outward currents at  $-60$  and  $-50$  mV, but the channels began to activate at  $-40$  mV. The activation time constants for individual subunits obtained by fitting the sigmoidal rising phase of the current traces with a Hodgkin-Huxley-related equation ( $m^4$ ) were  $0.72 \pm 0.06$  ms ( $n = 7$ ) for wild-type and  $0.77 \pm 0.09$  ms ( $n = 6$ ) for  $\Delta 2-40$  mutant channels at the standard test potential of  $+40$  mV. The two channel types showed the same voltage dependence of activation kinetics (e-fold enhancement per 21 mV depolarization; Fig. 2B). In order to compare the voltage dependence of activation of the channels, we fitted Boltzmann functions to the peak conductance–voltage relationships for wild-type and  $\Delta 2-40$  mutant channels (Fig. 2C). Fourth-order Boltzmann fits, which conform with the Hodgkin-Huxley-related equation used to approximate activation kinetics, yielded  $V_{1/2}$  values of  $-38.2 \pm 2.4$  mV

( $n = 8$ ) and  $-39.9 \pm 1.7$  mV ( $n = 6$ ) for Kv4.2 wild-type and  $\Delta 2-40$ , respectively, with slope factors of  $20.6 \pm 1.3$  and  $21.1 \pm 1.2$  mV. These  $V_{1/2}$  values for a single subunit represent the membrane potential for 6.25% of the maximal conductance and can serve as a rough estimate of the activation threshold, which was not significantly changed by the N-terminal truncation ( $P = 0.58$ ). Using the more conventional first-order Boltzmann fit we obtained  $V_{1/2}$  values of  $-1.2 \pm 2.2$  mV ( $n = 8$ ) for wild-type and  $-3.3 \pm 3.0$  mV ( $n = 6$ ) for  $\Delta 2-40$  mutant channels, indicating that the midpoints of their activation curves were not significantly different ( $P = 0.57$ ). The respective slope values for the fitted curves were  $16.2 \pm 0.9$  mV ( $n = 8$ ; wild-type) and  $16.3 \pm 0.8$  mV ( $n = 6$ ;  $\Delta 2-40$ ).

### Closed-state inactivation of Kv4.2 channels

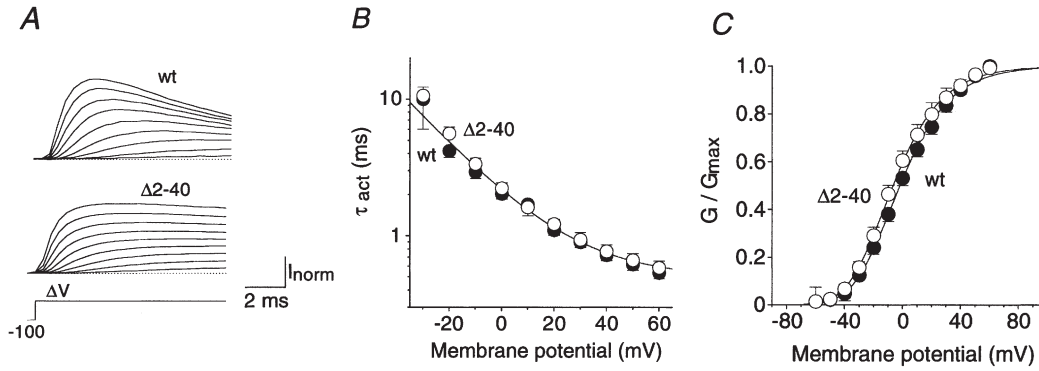
While in *Shaker* channels inactivation is tightly coupled to channel opening (Zagotta *et al.* 1989), *Shal* gene-related channels may readily inactivate from pre-open closed states as shown for Kv4.1 (Jerng *et al.* 1999). We intended to measure the kinetics of closed-state inactivation for both wild-type and  $\Delta 2-40$  mutant Kv4.2 channels. For this purpose we applied a double-pulse protocol with an intervening conditioning pulse of variable length to a potential below the activation threshold (e.g.  $-50$  mV; Fig. 3A). In such experiments, both the initial control pulse and the conditioning pulse were applied from a membrane potential of  $-100$  mV (Fig. 3A), where no steady-state inactivation was observed for Kv4.2 channels (see Fig. 4B). The current amplitudes evoked by the test pulse at the end of the protocol were normalized to control and plotted against the conditioning pulse duration (Fig. 3B).



**Figure 1. Inactivation kinetics of Kv4.2-mediated currents**

A, currents mediated by the wild-type (wt) and an N-terminal deletion mutant ( $\Delta 2-40$ ) of Kv4.2 expressed in HEK 293 cells and recorded in the whole-cell patch-clamp configuration during a 2.5 s voltage pulse to  $+40$  mV. Endogenous HEK 293  $K^+$  currents recorded following a prepulse to  $-35$  mV (not shown) have been subtracted from the total currents recorded after a prepulse to  $-100$  mV. Traces have been normalized to peak ( $I_{norm}$ ) and the dotted line represents zero current. Note the stronger Kv4.2  $\Delta 2-40$  inactivation, which causes a cross-over with the Kv4.2 wt trace. B, same recordings as in A on a 10-fold expanded vertical scale to illustrate intermediate and slow inactivation kinetics. Currents were fitted with three exponentials (continuous line). Inappropriate double-exponential fitting curves are shown as dashed lines. C, mean values for the fast ( $\tau_1$ , ■), intermediate ( $\tau_2$ , ▨) and slow components ( $\tau_3$ , □) of inactivation at  $+40$  mV for Kv4.2 wt ( $n = 13$ ) and  $\Delta 2-40$  mutant channels ( $n = 9$ ). The mean percentage of the total decay accounted for by each component is indicated.





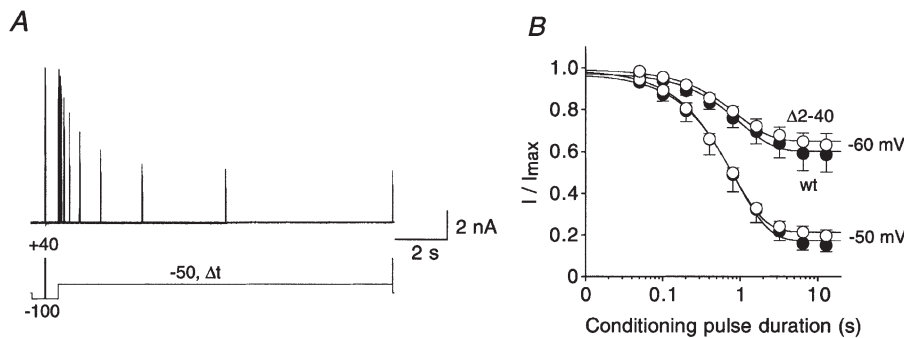
**Figure 2. Voltage dependence of Kv4.2 channel activation**

*A*, currents evoked by successive depolarizing steps from  $-100$  mV to potentials between  $-20$  and  $+60$  mV ( $\Delta V$ ). Note that activation kinetics are unaffected by the N-terminal deletion while inactivation proceeds more slowly for  $\Delta 2-40$  (lower traces) than for wt (upper traces). Dotted lines represent zero current. *B*, mean activation time constants ( $\tau_{act}$ ) for wt ( $\bullet$ ;  $n = 7$ ) and  $\Delta 2-40$  ( $\circ$ ;  $n = 6$ ) on a logarithmic scale plotted against the membrane potential applied during the test pulse. A single-exponential function was fitted to all plotted data points simultaneously. *C*, activation curves obtained with fourth-order Boltzmann fits to the normalized peak conductance values at different test potentials for wt ( $\bullet$ ) and  $\Delta 2-40$  ( $\circ$ ), respectively.

The results showed that Kv4.2 channels can enter inactivated states without having been activated. Single-exponential functions fitted to the data yielded time constants for the closed-state inactivation at  $-50$  mV of  $959 \pm 188$  ms ( $n = 6$ ) and  $774 \pm 79$  ms ( $n = 5$ ) for wild-type and  $\Delta 2-40$  mutant Kv4.2 channels, respectively. At  $-60$  mV closed-state inactivation occurred with a mean time constant of  $1.03 \pm 0.14$  s ( $n = 7$ ) for wild-type and  $912 \pm 138$  ms ( $n = 5$ ) for  $\Delta 2-40$  mutant channels. The slightly higher values obtained at  $-60$  mV may reflect the voltage dependence of closed-state inactivation, which was not further investigated, due to the small voltage range available; the inactivation observed at  $-70$  mV was usually too small to perform reliable kinetic analyses. At  $-60$  mV ( $P = 0.58$ ) and at  $-50$  mV ( $P = 0.64$ ), the small

differences between the mean closed-state inactivation kinetics for wild-type and  $\Delta 2-40$  mutant channels were not significant. Thus, N-terminal truncation of Kv4.2 channels seems to have significant effects on inactivation only if the open-state is involved in the inactivation reaction, as shown by the Kv4.2 A-type current decay kinetics in Fig. 1.

Figure 4*A* summarizes our kinetic analyses of outward current decay and closed-state inactivation over the voltage range studied, for both wild-type and  $\Delta 2-40$  mutant Kv4.2 channels. It demonstrates the slowing of both the fast ( $\tau_1$ ) and the intermediate component ( $\tau_2$ ) of A-type current decay by the N-terminal truncation, which apparently accelerates the slow component ( $\tau_3$ ).



**Figure 3. Onset of inactivation at potentials below activation threshold**

*A*, currents mediated by Kv4.2 wt channels recorded at  $+40$  mV using a double-pulse protocol with a conditioning pulse to  $-50$  mV of variable duration ( $\Delta t$ ). *B*, mean current amplitudes evoked by the second pulse in protocols as shown in *A*, relative to the amplitude obtained by the initial control pulse, were plotted against the duration of the conditioning pulse, on a logarithmic time base.  $\bullet$ , wt ( $n = 7$  for  $-60$  mV;  $n = 6$  for  $-50$  mV);  $\circ$ ,  $\Delta 2-40$  ( $n = 5$  for  $-60$  mV;  $n = 5$  for  $-50$  mV). Fitting curves represent single-exponential functions. Note that subtraction of endogenous HEK 293 currents from raw data such as shown in *A* resulted in a lower mean steady-state value at  $-50$  mV.

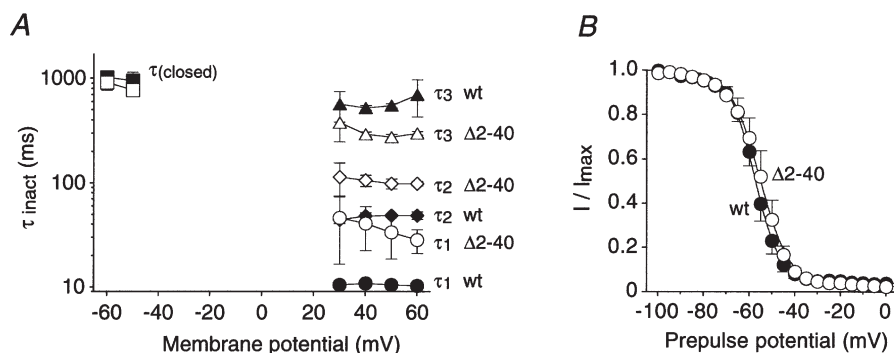
The voltage dependence of Kv4.2 steady-state inactivation was studied with a 5 s prepulse protocol using a test potential of +40 mV. The normalized results are shown in Fig. 4*B*. Boltzmann analysis yielded  $V_{1/2}$  values for the steady-state inactivation of  $-57.4 \pm 1.4$  mV ( $n = 6$ ) for wild-type and  $-54.5 \pm 2.2$  mV ( $n = 5$ ) for  $\Delta 2-40$  mutant channels. The  $V_{1/2}$  values were not significantly different ( $P = 0.29$ ). The slope factors were  $5.0 \pm 0.3$  mV for wild-type and  $5.3 \pm 0.7$  mV for  $\Delta 2-40$ . The finding that steady-state inactivation curves were almost identical for Kv4.2 wild-type and  $\Delta 2-40$  mutant channels, both above and below the activation threshold, supports the idea that the channels accumulate in an as yet undefined inactivated state, irrespective of the kinetics of inactivation transitions involving the open state. We suggest that the channels accumulate in the closed-inactivated state(s).

### Recovery of Kv4.2 channels from inactivation

One of the most striking features of *Shal* channels, which distinguishes them from *Shaker*-related channels, is their fast recovery from inactivation (Pak *et al.* 1991; Serodio *et al.* 1994). We investigated the involvement of transitions between closed-inactivated and pre-open closed state(s) in the pathway of recovery from inactivation. Thus, we characterized in detail the voltage dependence of the recovery kinetics for both wild-type and  $\Delta 2-40$  mutant Kv4.2 channels. In our double-pulse protocols (Fig. 5*A*) the membrane potential before the initial control pulse and during the recovery period was varied between  $-60$  (Fig. 5*A*) and  $-120$  mV (Fig. 5*B*). The recovery kinetics, obtained by plotting the relative amplitudes of the current activated by the second pulse against the duration of the interpulse interval, were always well described by a single-exponential function, although the onset of

inactivation occurred with multiple components (see Fig. 1). This is consistent with an accumulation of Kv4.2 channels in one inactivated state. The recovery from inactivation was strongly voltage dependent. This is illustrated in Fig. 5*C*, which shows, on a logarithmic time base, the results of our analysis of recovery from inactivation at  $-60$ ,  $-80$  and  $-120$  mV for Kv4.2 wild-type channels. The respective time constants were  $656 \pm 120$  ms ( $n = 4$ ),  $285 \pm 21$  ms ( $n = 4$ ) and  $69 \pm 11$  ms ( $n = 3$ ). The recovery time constants obtained at the same membrane potentials for  $\Delta 2-40$  mutant channels were not significantly different from wild-type ( $P$  values between 0.27 and 0.83; Fig. 5*D*), and the voltage dependence of the Kv4.2 recovery kinetics could be described by an exponential function suggesting an e-fold enhancement of recovery from inactivation per 24 mV hyperpolarization.

In our double-pulse protocols, a large fraction of the Kv4.2 channels passed through the open state before reaching the inactivated state(s), because the channels were quickly activated by the initial large depolarizing voltage jump. Since Kv4.2 channels also show transitions from the pre-open closed state(s) to the inactivated state(s), as shown in Fig. 3, we intended to measure the kinetics of recovery from such closed-state inactivation. For this purpose, after a short control pulse to +40 mV the channels were allowed to close and slowly reach steady-state inactivation during a 5 s conditioning pulse to  $-50$  mV. This was followed by a variable recovery period and a subsequent test pulse to +40 mV (Fig. 6*A*). The membrane potential before the initial pulse and during the recovery period was varied between  $-80$  (Fig. 6*A*) and  $-120$  mV (Fig. 6*B*). Kinetic analysis of the recovery from closed-state inactivation revealed no



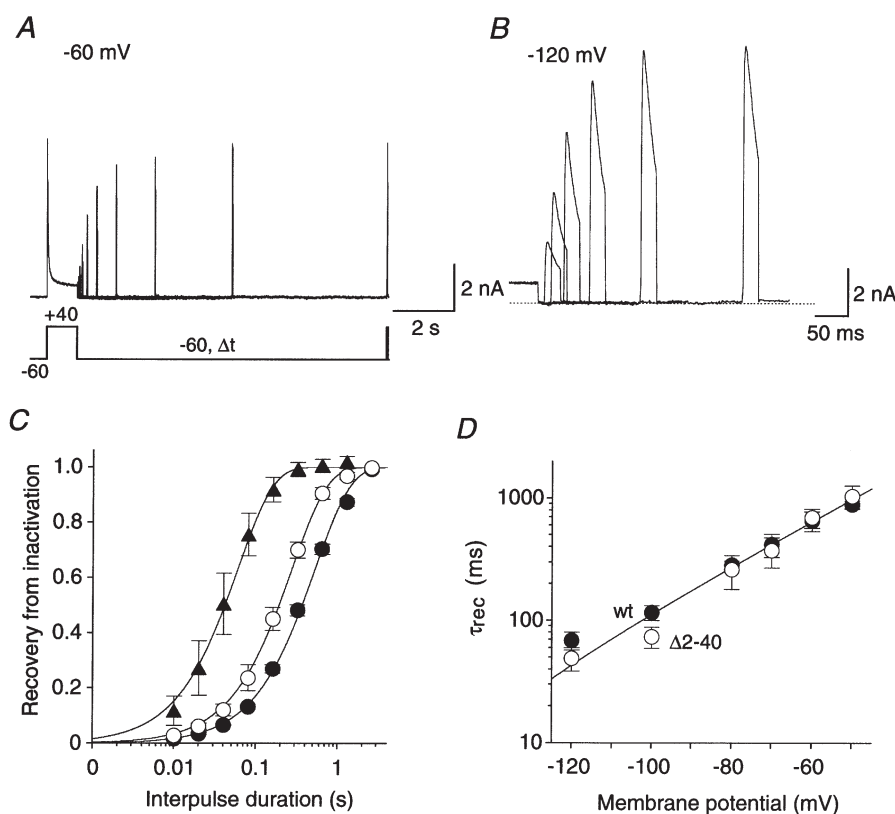
**Figure 4. Summary of kinetic and steady-state analysis of Kv4.2 channel inactivation**

*A*, mean inactivation time constants ( $\tau_{\text{inact}}$ ) measured at different membrane potentials for wt (filled symbols) and  $\Delta 2-40$  (open symbols). Circles, diamonds and triangles represent  $\tau_1$ ,  $\tau_2$  and  $\tau_3$ , respectively ( $n = 3$  for wt;  $n = 4$  for  $\Delta 2-40$ ), obtained from triple-exponential fits to the current traces obtained during depolarizing voltage jumps. Squares represent time constants of closed-state inactivation ( $\tau_{\text{closed}}$ );  $n = 5-7$ ). Lines connect individual data points for  $\tau_1$ ,  $\tau_2$ ,  $\tau_3$  and  $\tau_{\text{closed}}$ . Data are presented on a semi-log scale to better illustrate the differences between wt and  $\Delta 2-40$ . *B*, voltage dependence of steady-state inactivation. Normalized current values recorded after a 5 s conditioning pre-pulse were plotted against the pre-pulse potential for wt (●;  $n = 6$ ) and  $\Delta 2-40$  (○;  $n = 5$ ). Fitting curves represent first-order Boltzmann functions.

differences between wild-type and mutant Kv4.2 channels (Fig. 6C). The respective time constants obtained at a recovery potential of  $-80$  mV were  $241 \pm 19$  ms ( $n = 17$ ) and  $203 \pm 23$  ms ( $n = 15$ ;  $P = 0.21$ ). Both time constants were slightly smaller than, but not significantly different from, the mean values obtained for wild-type (285 ms;  $P = 0.30$ ) and mutant Kv4.2 channels (260 ms;  $P = 0.37$ ) using a conventional recovery protocol as shown in Fig. 5. Figure 6D summarizes the results of our kinetic analyses of recovery from closed-state inactivation for the recovery potentials  $-80$ ,  $-100$  and  $-120$  mV. The data suggest that the voltage dependence of recovery from closed-state inactivation is equal for wild-type and  $\Delta 2-40$  mutant channels. Progressively smaller deviations from the steady-state inactivation at  $-50$  mV prevented such analysis for recovery potentials positive to  $-80$  mV. However, the data points representing the time constants for  $-80$ ,  $-100$

and  $-120$  mV closely match the voltage dependence of recovery kinetics that was determined using conventional recovery protocols (compare Figs 5D and 6D). Taken together, our results show that wild-type and  $\Delta 2-40$  mutant Kv4.2 channels resemble each other with respect to all of the gating parameters tested so far, except inactivation transitions involving the open state, as revealed by the outward current decay kinetics.

Our recovery data presented in Figs 5 and 6 are consistent with the idea that Kv4.2 channels always accumulate in a closed-inactivated state, from which they directly recover, bypassing the open state. Prominent tail currents as observed for *Shaker* channels during recovery from inactivation at very negative membrane potentials and in the presence of high external  $K^+$  concentrations (Demo & Yellen, 1991; Ruppersberg *et al.* 1991) should therefore



**Figure 5.** Recovery of Kv4.2 channels from inactivation

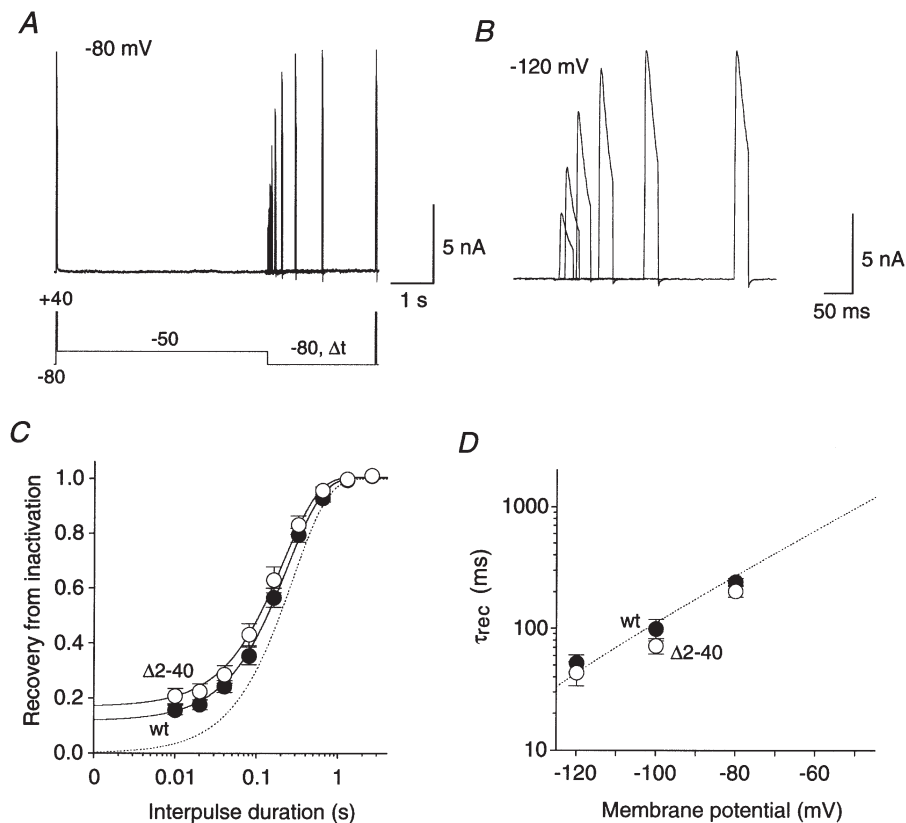
A, currents recorded with a double-pulse protocol at  $+40$  mV. Membrane potential before the initial voltage jump and during the variable interpulse interval ( $\Delta t$ ) was  $-60$  mV, where a large fraction of channels were already inactivated. B, currents recorded from the same cell with a protocol as in A but with a holding potential of  $-120$  mV, where steady-state inactivation was removed. Therefore, the maximum currents obtained at  $+40$  mV are larger than in A. Only the second current peak following various interpulse intervals is shown, and a different time scale was chosen. Note that recovery from inactivation was almost complete after 150 ms. Dotted line represents zero current. C, mean normalized currents recorded in wt channels during the second pulse plotted against the duration of the interpulse interval, on a logarithmic time base, for recovery potentials of  $-60$  mV (●;  $n = 4$ ),  $-80$  mV (○;  $n = 4$ ) and  $-120$  mV (▲;  $n = 3$ ). Single-exponential functions fitted to the data points yielded the recovery time constants. D, voltage dependence of recovery kinetics. A mono-exponential function was fitted simultaneously to all plotted data points representing the time constants of recovery ( $\tau_{rec}$ ) obtained for wt (●) and  $\Delta 2-40$  (○) on a logarithmic scale.

be absent in Kv4.2 channels. This assumption was tested by using protocols with voltage jumps from +40 or -50 mV to -120 mV, in symmetrical 130 mM K<sup>+</sup> (Fig. 7). As we expected, these experiments showed that progressive cumulative inactivation at +40 mV reduced the size of the initial tail currents during the recovery from inactivation at -120 mV (Fig. 7*A* and *B*). After 20 s of inactivation at +40 mV virtually no Kv4.2 channels resided in or revisited the open state during recovery from inactivation (Fig. 7*C*). Serving as a control experiment, the same result was obtained when the recovery from closed-state inactivation was examined at -120 mV in symmetrical 130 mM K<sup>+</sup> (Fig. 7*D*).

#### Kv4.2 wild-type and N-terminal deletion mutant differ in their deactivation kinetics

We measured the voltage dependence of the deactivation kinetics for Kv4.2 wild-type and  $\Delta 2-40$  mutant channels

using tail current protocols with test potentials between -110 and -50 mV following brief channel activation at +40 mV (Fig. 8*A*). The current decay due to channel deactivation could be fitted with a single-exponential function and the time constant at -110 mV for wild-type channels was  $1.6 \pm 0.1$  ms ( $n = 15$ ), significantly below the value for  $\Delta 2-40$  mutant channels ( $1.9 \pm 0.1$  ms;  $n = 15$ ;  $P < 0.05$ ). For both channel types the deactivation kinetics were slowed with increasing membrane depolarization, although with apparently different voltage dependence (Fig. 8*B*). As inferred from exponential fitting of the obtained data, the slowing was e-fold per 44 mV depolarization for wild-type and e-fold per 30 mV depolarization for  $\Delta 2-40$  mutant channels. Thus, the deactivation time constants measured at -50 mV, just below the activation threshold, differed most significantly ( $P < 0.01$ ) for wild-type and mutant Kv4.2 channels. Wild-type channel closing at this membrane potential



**Figure 6.** Recovery from closed-state inactivation

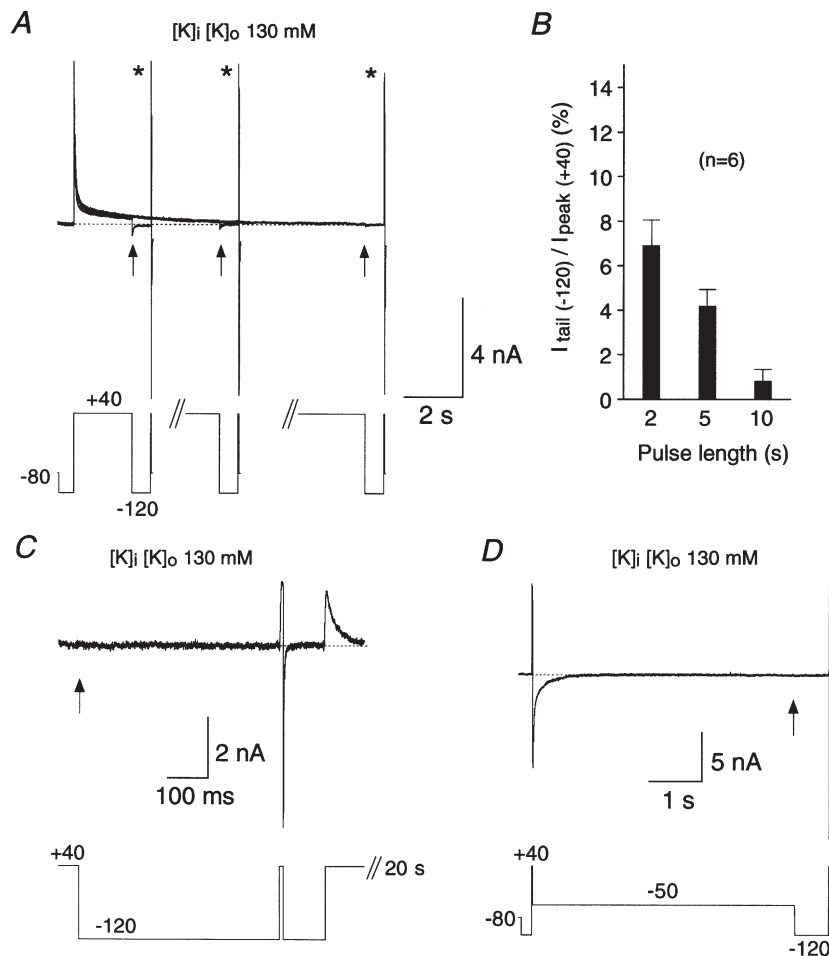
*A*, currents recorded at +40 mV with a specialized double-pulse protocol. Inactivation was induced at -50 mV and recovery was measured after variable times ( $\Delta t$ ) at -80 mV. *B*, currents recorded with the same protocol as in *A* but recovery at -120 mV was measured. Only the second current response is shown on a different time scale. *C*, kinetics of recovery from closed-state inactivation at -80 mV. Mean normalized current amplitudes were plotted against the interpulse duration, on a logarithmic time base, for wt (●;  $n = 17$ ) and  $\Delta 2-40$  (○;  $n = 15$ ). Fitting curves are based on single-exponential functions. The dotted line represents the fitting curve, which describes the recovery of wt channels at -80 mV from inactivation after channel opening, as shown in Fig. 5. *D*, mean time constants of recovery from closed-state inactivation for wt (●;  $n = 6-17$ ) and  $\Delta 2-40$  (○;  $n = 5-15$ ) plotted against the recovery potential. The voltage and time range shown (*x*- and *y*-axis, respectively) as well as the exponential fitting curve (dotted line) were adapted from Fig. 5*D*.



occurred with a mean time constant of  $5.9 \pm 0.3$  ms ( $n = 15$ ) while the respective value for  $\Delta 2-40$  mutant channels was  $11.9 \pm 0.6$  ms ( $n = 15$ ). Our results suggest that the N-terminal truncation stabilizes the open-state of Kv4.2 channels, which causes both slowed deactivation and slowed inactivation.

The closing of  $K^+$  channels can be influenced in a voltage-dependent manner by the use of ions with different permeation properties, like  $Rb^+$  and  $K^+$  (Sala & Matteson, 1991).  $Rb^+$  ions can slow the deactivation process due to their higher pore occupancy. In order to study the

permeant ion dependence of both deactivation and inactivation we recorded inward tail currents and transient outward currents, respectively, mediated by Kv4.2 wild-type channels in symmetrical 130 mM  $K^+$  and symmetrical 130 mM  $Rb^+$  (Fig. 9A and C). We used the same tail current protocol as under standard conditions. The test potential was, however, varied only between  $-110$  and  $-60$  mV, because the  $Rb^+$  ions caused a hyperpolarizing shift in the conductance–voltage relationship (not shown). The measured deactivation kinetics and their voltage dependence were strongly



**Figure 7.** No transient occupancy of open state during recovery after accumulation in closed-inactivated state(s)

Recovery of Kv4.2 wt channels measured at  $-120$  mV and in symmetrical 130 mM  $K^+$  to test whether channels visit the open state during recovery from inactivation. *A*, depolarizing pulses to  $+40$  mV of variable duration (2, 5 and 10 s) were followed by negative voltage jumps to  $-120$  mV. Three current traces recorded from the same cell are superimposed. Note that tail currents (arrows) vanish as cumulative inactivation proceeds. In all cases recovery from inactivation was complete after 640 ms, as indicated by the brief control outward currents (asterisks) followed by deactivating inward currents at the end of each pulse protocol. *B*, percentage of tail current amplitude (arrows in *A*) relative to the control outward current amplitude (asterisks in *A*) at the different pulse durations. Results from 6 cells are summarized. *C*, the holding potential was manually set to  $+40$  mV, and currents were recorded every 20 s at high resolution ( $50 \mu s$  sample interval). Tail currents were absent at the beginning of a 640 ms pulse to  $-120$  mV (arrow), but recovery during this period was complete. *D*, protocol similar to the ones used in Fig. 6, but with a fixed recovery time (640 ms) and fixed recovery potential ( $-120$  mV). No tail currents (arrow) but complete recovery could be observed. The dotted lines in *A*, *C* and *D* represent zero current.

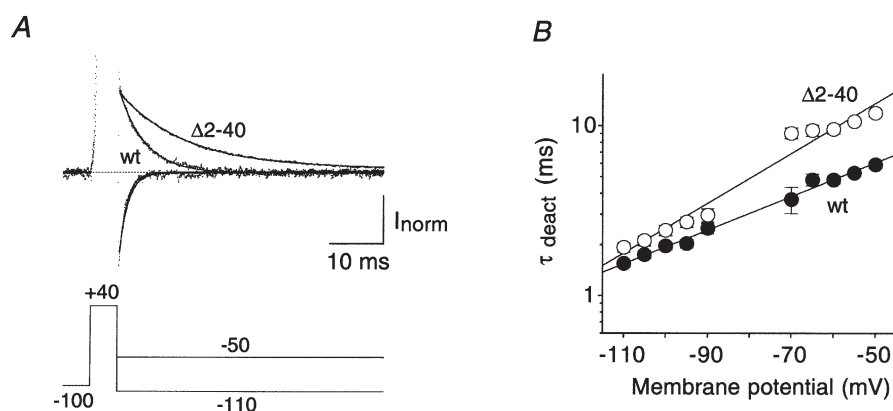
influenced by the permeant ion species. At  $-110$  mV channel closing occurred with mean time constants of  $1.3 \pm 0.2$  ms ( $n = 7$ ) and  $4.6 \pm 0.7$  ms ( $n = 10$ ) in symmetrical  $K^+$  and symmetrical  $Rb^+$ , respectively. Increasing depolarization slowed the deactivation kinetics, however, with slightly different voltage dependence for the two ion species (Fig. 9B). In symmetrical  $K^+$  an e-fold slowing per 54 mV depolarization led to a mean deactivation time constant of  $3.4 \pm 0.3$  ms at  $-60$  mV, while in symmetrical  $Rb^+$  a time constant of  $13.6 \pm 2.0$  ms was reached, due to an e-fold slowing per 46 mV depolarization. Next, we measured the inactivation kinetics of outward currents at  $+40$  mV in symmetrical  $K^+$  and symmetrical  $Rb^+$  (Fig. 9C). The inactivation kinetics obtained under symmetrical ionic conditions are summarized in Fig. 9D. As indicated in Fig. 9C, the increase in the external  $K^+$  concentration from 5 mM (standard conditions) to 130 mM caused a significant decrease in the mean time constants for fast ( $P < 0.01$ ) and intermediate ( $P < 0.01$ ) inactivation ( $5.8 \pm 0.5$  and  $36.5 \pm 2.6$  ms, respectively, accounting for  $63.1 \pm 2.8$  and  $31.9 \pm 2.2\%$  of the total decay;  $n = 5$ ). The slow time constant in symmetrical 130 mM  $K^+$  was  $588 \pm 72$  ms ( $n = 5$ ), not significantly different from that in standard conditions ( $P = 0.36$ ), and accounted for  $5.0 \pm 0.9\%$  of the total decay. In symmetrical 130 mM  $Rb^+$  the mean ( $n = 6$ ) fast time constant was  $13.9 \pm 2.0$  ms ( $66.9 \pm 3.4\%$ ), significantly slower than in symmetrical  $K^+$  ( $P < 0.05$ ). Similarly, the intermediate time constant was significantly slower in symmetrical  $Rb^+$  ( $65.8 \pm 9.4$  ms;  $26.0 \pm 3.1\%$ ;  $n = 6$ ) than in symmetrical  $K^+$  ( $P < 0.05$ ); however, the mean slow time constant was not significantly changed ( $693 \pm 121$  ms;  $7.1 \pm 0.5\%$ ;  $n = 6$ ) by  $Rb^+$  ( $P = 0.48$ ). In summary, the high concentration of  $Rb^+$  ions caused a

slowing of Kv4.2 channel deactivation and a slowing of the fast and intermediate components of Kv4.2 channel inactivation. These results are in accordance with the idea that channel closing, i.e. the occupancy of a pre-open closed state, plays a significant role in Kv4.2 channel inactivation gating.

### Simulating Kv4.2 inactivation gating

Previously, Jerng and coworkers (Jerng *et al.* 1999) developed a model for Kv4.1 channel gating (Model 1; Fig. 10A), in which channels become inactivated if they visit an open-inactivated state ( $I_o$ ), a closed-inactivated state ( $I_c$ ), or a deep-inactivated state ( $I_D$ ). The two main kinetic features of Model 1 are, first, that Kv4.1 channels have to return from the short-lived open-inactivated state  $I_o$  via the open state to the pre-open closed state before they can enter the closed-inactivated state,  $I_c$ , and second, that cumulative Kv4.1 channel inactivation is a two-step reaction involving the closed-inactivated state,  $I_c$ , and the deep-inactivated state,  $I_D$ .

We developed a Markov state model to simulate the Kv4.2 channel currents recorded for this study (see Methods). We started with the gating scheme that Jerng and coworkers (Jerng *et al.* 1999) used to fit their Kv4.1 data (Model 1; Fig. 10A), and adjusted  $\alpha_0$  and  $\beta_0$ , in order to match the simulated activation characteristics ( $V_{1/2} = -1.9$  mV;  $k = 16.5$  mV; not shown) to our experimentally obtained  $G-V$  curves. Furthermore, we used different inactivation rates (Table 1) in order to obtain current decay kinetics that resembled more closely the experimental Kv4.2 data. Figure 10 shows the simulated currents based on voltage protocols from this study (Figs 1A and 5B), which revealed key kinetic features of Kv4.2 channel



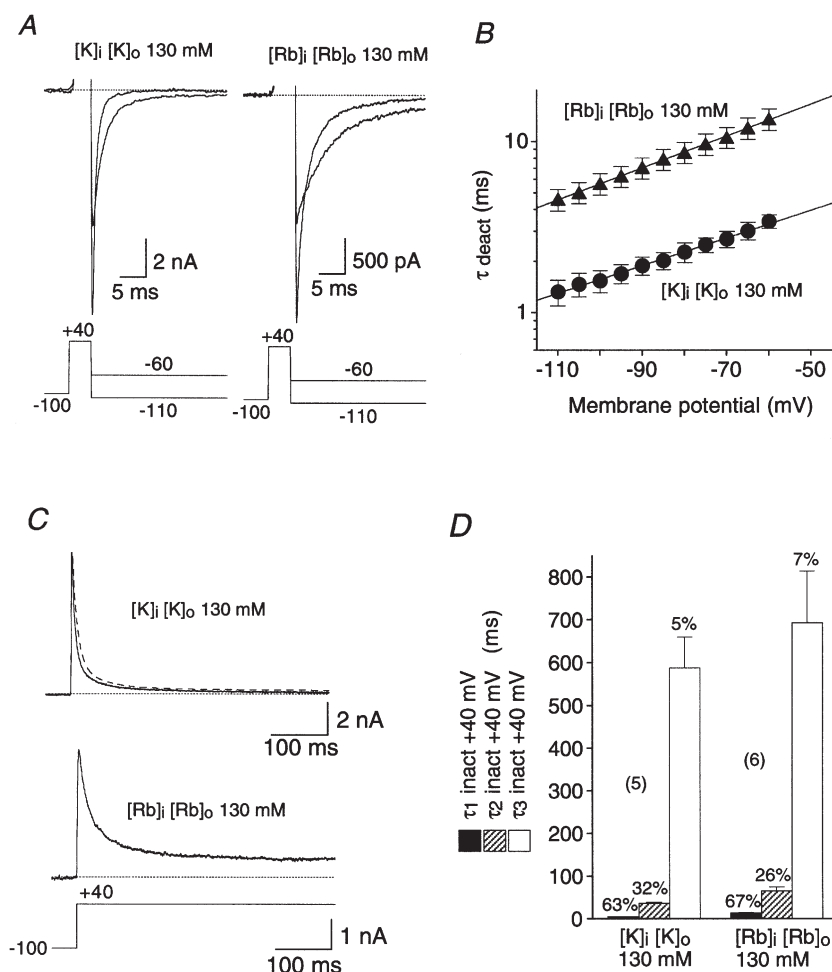
**Figure 8.** Deactivation kinetics of Kv4.2 channels

A, tail currents 5 ms after channel activation recorded during voltage jumps from  $+40$  to  $-110$  mV (fast decaying inward currents) and from  $+40$  to  $-50$  mV (slowly decaying outward currents) for wt and  $\Delta 2-40$ . Dotted line represents zero current and outward currents during the activation phase are truncated. Single-exponential functions were fitted to the current relaxations. Note the large difference in deactivation kinetics between wt and  $\Delta 2-40$  at  $-50$  mV, which was less obvious at  $-110$  mV. B, mean time constants of deactivation ( $\tau_{\text{deact}}$ ) for wt (●) and  $\Delta 2-40$  (○) plotted on a logarithmic scale at different membrane potentials ( $n = 15$  for both channel types). The voltage dependence of deactivation kinetics was approximated by mono-exponential functions.

inactivation. The currents simulated with Model 1 and the parameter values listed in Table 1 exhibited multi-exponential inactivation kinetics during a depolarizing pulse to +40 mV, very similar to those obtained experimentally (Fig. 10B). However, our simulations failed to reproduce the experimentally observed fast recovery from inactivation at -120 mV. This indicated that Model 1 was inappropriate (Fig. 10C).

We were particularly interested in the path Kv4.2 channels take to reach an inactivated state and their

distribution among the different inactivated states at a given time point. To address this question, we calculated the relative occupancy of each state of the Markov chain model during the voltage pulse as a function of time (fractional occupancy time course; Fig. 10D). For ease of display, we plotted only the pre-open closed ( $C_4$ ), open (O) and inactivated (I) states. The fractional occupancy timecourse shows that in Model 1 channels exit both  $I_O$  and  $I_C$ , and by the end of the pulse all channels reside in the  $I_D$  state. Since exit from  $I_D$  is slow, the time course for recovery from inactivation is very slow too.



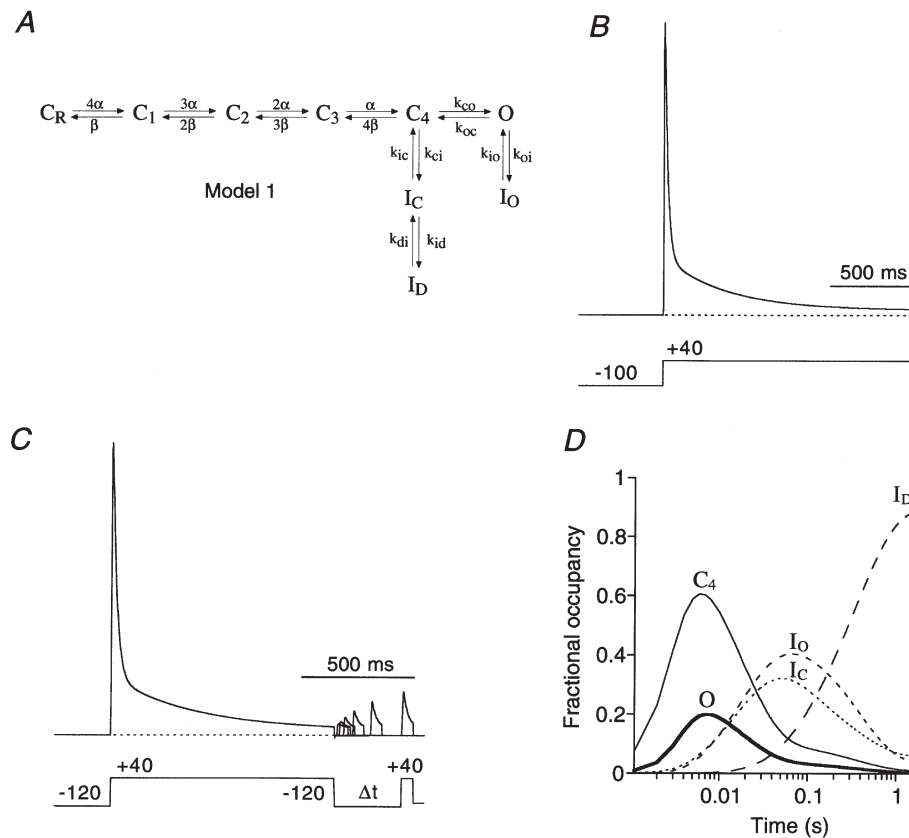
**Figure 9.** Deactivation and inactivation of Kv4.2 wild-type channels under varied ionic conditions

*A*, inward tail currents measured with the same protocol as shown in Fig. 8A but in 130 mM symmetrical  $K^+$  (left) and in 130 mM symmetrical  $Rb^+$  (right). Dotted lines represent zero current and outward currents during the activation phase are truncated. *B*, mean deactivation time constants obtained from single-exponential fits to the tail currents, plotted on a logarithmic scale against the membrane potential for symmetrical  $K^+$  (●;  $n = 7$ ) and symmetrical  $Rb^+$  (▲;  $n = 10$ ). Mono-exponential functions describe the voltage dependence of the deactivation kinetics. *C*, outward currents carried by  $K^+$  (upper trace) and  $Rb^+$  (lower trace) under symmetrical ionic conditions were recorded during depolarizing voltage jumps from -100 to +40 mV. The upper trace was obtained from the Kv4.2 wt-expressing cell shown in Fig. 1. The triple-exponential fit from Fig. 1B was normalized and superimposed on the current recorded in symmetrical 130 mM  $K^+$  (dashed line). Dotted lines represent zero current. *D*, mean values for the fast ( $\tau_1$ , ■), intermediate ( $\tau_2$ , ▨) and slow components ( $\tau_3$ , □) of inactivation at +40 mV for symmetrical  $K^+$  ( $n = 5$ ) and symmetrical  $Rb^+$  ( $n = 6$ ). The mean percentage of the total decay accounted for by each component is indicated.

Further simulations based on Model 1 with different entry and exit rates for inactivated states failed to adequately describe Kv4.2 channel gating. The main criteria for an adequate simulation were a fast time course of recovery from inactivation at hyperpolarized potentials, and a pronounced and multi-exponential inactivation during strong depolarization. This could only be obtained when we introduced two major alterations to Model 1. First, we removed the deep-inactivated state,  $I_D$ , since accumulation in  $I_D$  was responsible for the extremely slow recovery from inactivation. Instead, we introduced an additional inactivated state,  $I_S$ , which can be entered from  $I_O$  to maintain the experimentally observed three components of current decay in the simulations. Second, we connected all closed states to inactivated states, in analogy to a Kv4.3 gating model used by Greenstein *et al.*

(2000). We refer to the new model (Model 2; Fig. 11A) as an ‘allosteric model of inactivation’, as suggested by Beck *et al.* (2001). In such an allosteric model the ‘allosteric factor’  $f$  (Table 2), similar to the scaling factor used by Greenstein *et al.* (2000), defines the coupling of closed-state inactivation to the individual transition steps along the activation pathway ( $C_R$  to  $C_4$  in Model 2). Thus, inactivation is coupled to voltage-dependent activation. The possibility that  $I_S$  may connect  $I_O$  and  $I_{C_4}$ , as indicated in Model 2 (Fig. 11A), was not tested in this study (but see Discussion).

The zero voltage rate term  $\alpha_0$  for the forward direction in the closed states was adjusted (Table 2) to make the allosteric model simulate an activation behaviour ( $V_{1/2} = -1.1$  mV;  $k = 17.4$  mV) similar to that obtained experimentally for Kv4.2 channels. In order to fit the



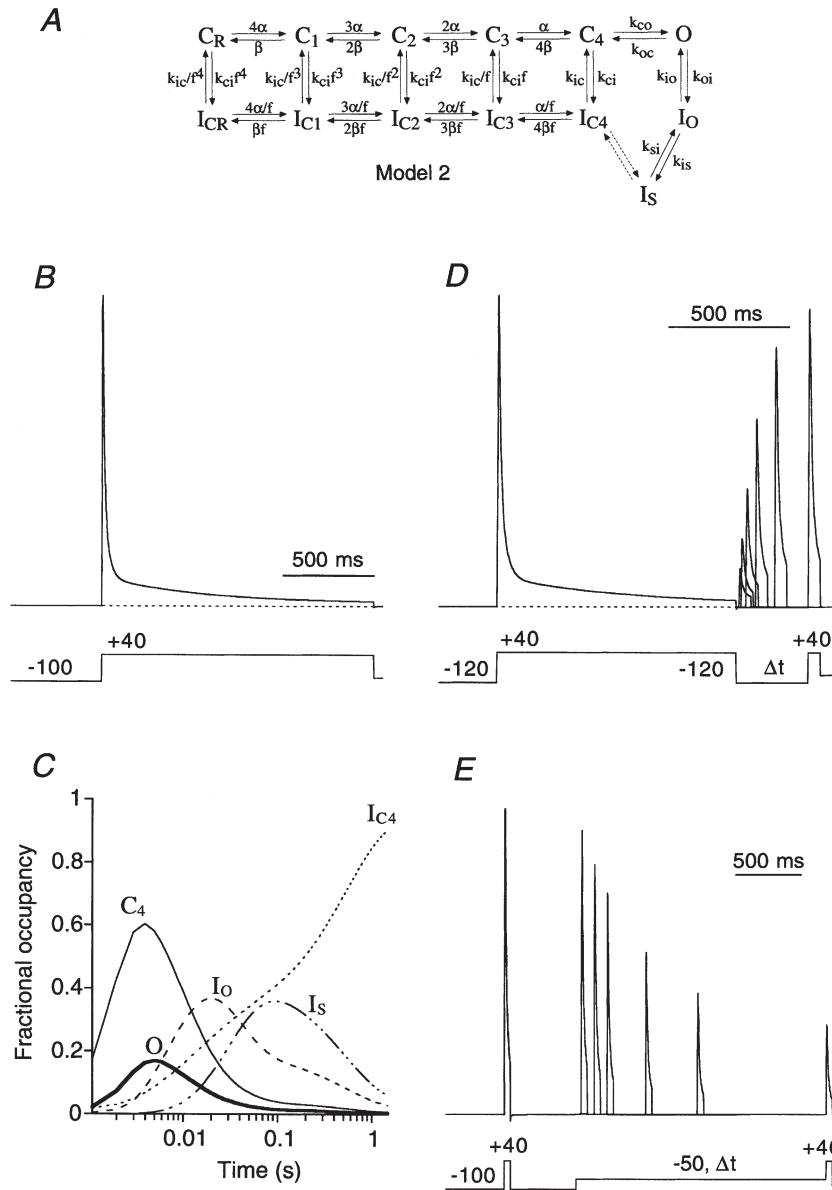
**Figure 10. Predictions of a Kv4.1 kinetic model with adjusted parameters for Kv4.2**

A, the  $K^+$  channel state diagram (Model 1) used for these simulations includes inactivation both from a pre-open closed state ( $C_4 \rightarrow I_C$ ) and from the open state ( $O \rightarrow I_O$ ). The closed-state inactivation also accesses a deeper inactivated state ( $I_D$ ). Transitions between states are represented by arrows; the values of the simulation parameters are shown in Table 1. B, simulation, with a pulse protocol like the one used in Fig. 1, to assess the rate of channel inactivation. Entry into the inactivated states has a fast, intermediate and slow component and inactivation is largely complete for a 1.5 s depolarizing pulse to +40 mV (dotted line represents zero current). C, simulation of the recovery from inactivation. Model 1 and parameters in Table 1 were used with a recovery protocol as in Fig. 5B. The model fails to simulate experimental data because recovery is very slow, even at -120 mV. D, the fractional occupancy of Model 1 gating states is plotted as a function of the inactivation pulse time on a logarithmic time base. For clarity, only the occupancy of the pre-open closed state ( $C_4$ ), the open state (O) and the three inactivated states ( $I_C$ ,  $I_O$  and  $I_D$ ) are shown. During the 1.5 s depolarizing pulse to +40 mV, channels largely accumulate in the deep inactivated state,  $I_D$ , where they get trapped.



model simulations to our data, the rate of recovery from inactivation was determined by the rate constant  $k_{ic}$ . The fast decay component of simulated current traces showed a high sensitivity to  $k_{oi}$ , and while  $k_{ei}$  played the major

role in setting the intermediate time constant of inactivation,  $k_{si}$  and  $k_{is}$  determined the ultra slow component of inactivation. Thus, each of the components of the inactivation kinetics were found to be controlled by



**Figure 11. Predictions of an allosteric model for Kv4.2 channel inactivation**

*A*, new Kv4 state diagram (Model 2), in which channels can inactivate from all closed states along the activation pathway. The allosteric factor  $f$  defines the coupling of closed-state inactivation and the voltage-dependent activation transitions leading from  $C_R$  to  $C_4$ ; arrows reflect the transitions between states. Values for the simulation parameters are given in Table 2. *B*, simulation to assess the rate of channel inactivation with a pulse protocol like the one used in Fig. 1. Entry into the inactivated states has a predominant fast component followed by a smaller intermediate and an even smaller slow component. Inactivation is largely complete for a 1.5 s pulse to +40 mV (dotted line represents zero current). *C*, the fractional occupancy of Model 2 gating states is plotted as a function of the inactivation pulse time on a logarithmic time base. For clarity, only the occupancy of the pre-open closed state ( $C_4$ ), the open state ( $O$ ) and the three inactivated states ( $I_{C4}$ ,  $I_O$  and  $I_S$ ) are shown. By the end of the 1.5 s depolarizing pulse to +40 mV a large fraction of channels has accumulated in  $I_{C4}$ . *D*, simulation of fast recovery from inactivation at -120 mV. *E*, simulation of the entry into closed-state inactivation with a pre-pulse to -50 mV, which did not activate a significant number of channels.

one or two crucial parameters in the Markov state scheme. Simulations based on Model 2 with the parameter values listed in Table 2 adequately described the Kv4.2 experimental data. The currents decayed with one fast ( $\tau = 9$  ms; 61%), one intermediate ( $\tau = 27$  ms; 31%) and one slow component ( $\tau = 544$  ms; 9%), like the measured Kv4.2 wild-type currents in response to a long depolarizing pulse (Fig. 11B). Since in the allosteric model the opening step is reverse-biased (Table 2) and inactivation from the open state is relatively unstable (Table 2), channels mainly inactivate from the pre-open closed state,  $C_4$ , and accumulate in  $I_{C4}$  during strong depolarization. Plotting the fractional occupancy time course (Fig. 11C) showed that channels start to enter  $I_{C4}$  early during the depolarizing pulse. After opening, the channels also occupy  $I_0$  and  $I_s$ , which slows the entry into  $I_{C4}$ . Finally, the channels slowly exit  $I_0$  and  $I_s$ , close, and accumulate in  $I_{C4}$ , from where the majority of channels directly recover, bypassing the open state, as soon as the membrane polarity is reversed. The simulated time course of recovery from inactivation (Fig. 11D;  $\tau = 79$  ms at  $-120$  mV) closely resembled that measured for Kv4.2 currents. The model was also able to adequately simulate the onset of closed-state inactivation at  $-50$  mV (Fig. 11E), albeit with slightly faster kinetics ( $\tau = 521$  ms) as seen on average experimentally.

## DISCUSSION

The main conclusion that can be drawn from the results of our experiments is that in response to membrane depolarizations of variable strength, Kv4.2 channels always accumulate in the closed-inactivated state(s), from which they directly recover, bypassing the open state. The simulations done with an allosteric model nicely reproduced the measured Kv4.2 gating kinetics, especially the rapid time course of recovery from inactivation.

### Inactivation reactions underlying A-type current decay

We compared our results with kinetic analyses done previously on Kv4.2 channels transiently expressed in HEK 293 cells (Perez-Garcia *et al.* 1999; Zhu *et al.* 1999). In contrast to previous studies, we have determined endogenous HEK 293 outward currents, which may contaminate decay kinetics. Thus, their subtraction is essential for a detailed kinetic analysis of macroscopic Kv4.2 currents. Furthermore, we utilized very long depolarizing pulses, in order to also cover extremely slow components of outward current decay. The results revealed triple-exponential decay kinetics for Kv4.2-mediated currents. Similar multi-exponential decay kinetics have been described for Kv4.1 channels (Jerng & Covarrubias, 1997; Jerng *et al.* 1999). However, Kv4.1 and Kv4.2 current inactivation kinetics are distinct. Kv4.1-mediated currents show only a small fast component, whereas in the Kv4.2 current kinetics the fast component accounts for the largest fraction of the

total decay. This quantitative difference may be the reason why the slowing of fast and intermediate decay components produced by N-terminal truncation can be seen in Kv4.2-mediated currents, while similar deletion mutants of Kv4.1 apparently lack the fast component of inactivation (Jerng & Covarrubias, 1997). A complete removal of rapid inactivation after deletion of N-terminal domains can be observed in *Shaker* B channels (Hoshi *et al.* 1990); however, Kv4.1 channels do not undergo classical N-type inactivation since Kv4.1 inactivation kinetics are insensitive to internal tetraethylammonium (Jerng & Covarrubias, 1997), an open channel blocker, which interferes with *Shaker* N-type inactivation (Choi *et al.* 1991). We have not performed comparable pharmacological experiments; however, the distinct effects of the 40 amino acid N-terminal truncation observed in our experiments (slowing of both fast and intermediate components, rather than complete removal of the fast component of current decay) clearly show that inactivation of Kv4.2 channels is different from the N-type inactivation of *Shaker*-related channels.

Jerng & Covarrubias (1997) have also shown that external tetraethylammonium, which interferes with *Shaker* channel C-type inactivation (Choi *et al.* 1991), does not affect the Kv4.1 current decay kinetics, which are therefore unlikely to represent C-type inactivation. Similarly, the results we obtained in experiments with high  $K^+$  concentrations suggest that the slow decay phase of Kv4.2-mediated currents is distinct from *Shaker*-like C-type inactivation. The latter is slowed by increasing the external  $K^+$  concentration (Baukowitz & Yellen, 1995), whereas the slow Kv4.2 inactivation component was not affected by an increase in the external  $K^+$  concentration (compare Figs 1C and 9D). We did, however, observe an enhancement of fast Kv4.2 inactivation in symmetrical  $K^+$ , which nicely compares to the effect of elevated external  $K^+$  concentration described for somatodendritic A-type currents recorded from isolated hippocampal neurons (Kirichok *et al.* 1998). The recovery of C-type-inactivated *Shaker* channels tends to be slow, occurring in the time range of many seconds, but Kv4.2 channels completely recovered from inactivation within a few hundred milliseconds (see Fig. 5). We did not study recovery from inactivation in high external  $K^+$ , but Jerng & Covarrubias (1997) have shown that elevated  $K^+$  concentrations cause a slowing of Kv4.1 channel recovery, which is in contrast to the recovery-enhancing effect of high external  $K^+$  concentrations observed for *Shaker*-related channels (Demo & Yellen, 1991). It is obvious that Kv4 channel inactivation does not follow classical *Shaker* N- and C-type inactivation.

### Simulations

According to Model 1 (Fig. 10A) the Kv4 current decay is a consequence of the entry into a short-lived open-inactivated state ( $O \rightarrow I_0$ ), the accumulation in a closed-inactivated state ( $I_C$ ), and the final trapping in a deep-

inactivated state ( $I_D$ ). The model simulated adequately the activation, inactivation and deactivation kinetics of macroscopic Kv4.1 currents (Jerng *et al.* 1999). The final trapping in a deep-inactivated state is, however, not compatible with fast recovery from inactivation (see Fig. 10C), as was experimentally observed for both Kv4.1 and Kv4.2 channels. Therefore, we used an alternative gating model (Model 2 in Fig. 11A), which was derived from Model 1 by modifying the entry and exit rates for the two inactivated states  $I_O$  and  $I_C$ , by deleting  $I_D$  and instead coupling  $I_O$  to another inactivated state,  $I_S$ , and finally by connecting all closed states along the activation pathway to inactivated states (allosteric model; Beck *et al.* 2001). Simulations with this model produced a fast outward current component, which accounted for a large fraction of the total decay, and both recovery from inactivation and closed-state inactivation could be simulated adequately.

### Involvement of channel closing in Kv4 inactivation gating

Both in Model 1 and in Model 2, the prerequisite for accumulation in the closed-inactivated state(s) is a transition from  $I_O$  back to the open state, and the subsequent closing of channels, because  $I_O$  and  $I_C$  are not connected. In fact, when slowing the channel closing rate by a factor of 20 in Model 1, Jerng and coworkers were able to reproduce the macroscopic inactivation time course that they found in slowly deactivating Kv4.1 mutants (Jerng *et al.* 1999). The prominent feature of the mutant Kv4.1 inactivation was a slowing of the intermediate and slow components of decay (Jerng *et al.* 1999). In the Kv4.2  $\Delta 2-40$  N-terminal deletion mutant, which also exhibited slowed deactivation kinetics, we observed on average a 3-fold slowing of the fast component, a 2-fold slowing of the intermediate component, and a 2-fold enhancement of the slow component of current decay (see Fig. 1C). This is in contrast to the slowly deactivating Kv4.1 mutant studied by Jerng *et al.* (1999). Also, the effect of the 40 amino acid deletion on Kv4.2 current inactivation could not be exactly reproduced when  $K^+$  was replaced by  $Rb^+$ , which, due to its higher pore occupancy, slows deactivation (Sala & Matteson, 1991). In symmetrical  $Rb^+$  both the fast and the intermediate components of Kv4.2 current decay were slowed on average 2-fold relative to symmetrical  $K^+$ , while the slow component was unaffected by the permeant ion species (see Fig. 9D). Apparently, the effect of N-terminal truncation on Kv4.2 gating kinetics is distinct, leading to a stabilization of the open state, which causes, independently of each other, a pronounced slowing of deactivation at moderately depolarized potentials and a slowing of inactivation transitions involving the open state.

Interestingly, slowing the open-state inactivation rate ( $O \rightarrow I_O$ ) in Model 1 and Model 2 produces single-exponential tail current kinetics, while rapidly inactivating models always deactivate with double-exponential kinetics (not

shown). A prominent second component of tail current decay, probably representing channels that exit  $I_O$  via the open state before they close, was observed for Kv4.1 channels (Jerng *et al.* 1999), but not in our experiments (see Fig. 8A), although Kv4.2 channels exhibit a fairly rapid initial component of inactivation. Therefore we envisage the possibility that open-inactivated Kv4.2 channels can close. This would be in agreement with models that have been used previously to model A-type channel gating (Solc & Aldrich, 1990; Campbell *et al.* 1993; Greenstein *et al.* 2000). Kv4 channel closing may involve one or more subconductance states, which correspond to  $I_S$ . Interestingly, Campbell *et al.* (1993) adapted their model from one that was originally developed for voltage-dependent  $Na^+$  channels (Patlak, 1991). It has been shown that  $Na^+$  channels can inactivate from a closed state (Horn *et al.* 1981), and that they deactivate before they recover from inactivation (Kuo & Bean, 1994).

An important feature of both Model 1 and Model 2 is the reverse-biased opening step ( $k_{oc} > k_{co}$ ; see Tables 1 and 2 and Jerng *et al.* 1999). It allows cumulative inactivation in  $I_{C4}$  even at strongly positive membrane potentials. Such an absorbing inactivation transition, leading from  $C_4$  to  $I_{C4}$ , implies a low peak open probability (see Fig. 10D and 11C). It is noteworthy that Solc & Aldrich (1990), who studied native *Drosophila non-Shaker* A-type channels, disfavoured a model with an inactivation transition occurring from the final pre-open closed state, because it produced a lower maximal open probability than was observed experimentally. Clarification of this question and the existence of a putative subconductance state(s) has to await a detailed analysis of single-channel properties.

### Closed-inactivated state(s)

The recovery of  $Na^+$  channels from inactivation is fast and voltage dependent, because it occurs from a closed-inactivated state. Here we show that the same is true for Kv4.2 channels. Our experimental evidence, that Kv4.2 channels accumulate in and directly recover from a closed-inactivated state bypassing the open state, can be summarized as follows. First, wild-type and  $\Delta 2-40$  mutant Kv4.2 channels recover from inactivation with identical kinetics (Fig. 5D), although their open-state inactivation kinetics are different (Fig. 1A). Second, for the two channel types, the recovery kinetics remain the same, no matter whether the channels opened ( $I_O \rightleftharpoons O \rightarrow C_4 \rightarrow I_{C4}$ ; Fig. 5) or not ( $C_4 \rightarrow I_{C4}$ ; Fig. 6) before they inactivated. Third, the recovery kinetics are strongly voltage dependent (Fig. 5C), the voltage dependence is identical for wild-type and  $\Delta 2-40$  mutant channels (Fig. 5D), and it is not changed by channel opening prior to inactivation (Fig. 6D). Finally, if cumulative inactivation of Kv4.2 channels in  $I_{C4}$  is complete, no inward tail currents can be observed, even in the presence of high external  $K^+$  concentrations, during recovery at very negative

membrane potentials (Fig. 7). These findings exclude a prominent involvement of the open state in the recovery pathway. By contrast, *Shaker* channels inactivate exclusively via the open state and have to release the inactivation ball before they can close, as indicated by slow tail currents during the recovery from inactivation (Demo & Yellen, 1991; Ruppersberg *et al.* 1991). However, inactivation from a closed state has been described for the I2 mutant of *Shaker* B channels (Ayer & Sigworth, 1997). The mutation (L382I) was introduced immediately adjacent to the S4–S5 linker at the end of the S4 region. It makes *Shaker* B channels behave more like voltage-dependent Na<sup>+</sup> channels and, according to this study, like Kv4.2 channels. This raises the question whether only small structural changes may underlie the distinct gating behaviours of *Shaker*- and *Shal*-related channels.

### Concluding remarks

The predominant mechanism of Kv4.2 channel inactivation from a pre-open closed state is correlated with their important physiological role in dendritic signal propagation. Closed-state inactivation of somato-dendritic A-type channels can lead to spike amplification since these inactivated states can be reached as a consequence of subthreshold EPSPs in dendrites (Magee & Johnston, 1997). Kv4.2 inactivation transitions involving the open state are equally important in order to counterbalance sudden extensive depolarizations. As our experiments showed, recovery will be fast irrespective of the initial inactivation reaction ( $C_4 \rightarrow I_{C4}$  or  $O \rightarrow I_0$ ), which guarantees a high A-type channel availability. In fact, a family of Kv channel-interacting proteins (KChIPs) has recently been described (An *et al.* 2000), which are associated with native Kv4 channels. KChIPs dramatically increase the functional expression of Kv4 channels and markedly accelerate their recovery from inactivation. A detailed kinetic analysis of Kv4 channel gating in the presence of KChIPs will shed further light on the inactivation mechanism of these channels and the structures involved.

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